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# Top-Quark Decays and QCD corrections

QCD@Colliders

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# Outline

- SM Top Decays
- QCD corrections
  - $O(\alpha_s)$
  - $O(\alpha_s^2)$
- Bonus material: “Unconventional decays”



# SM Top Decays

Predominantly:  $t \rightarrow W^+ b$

- $W^+ \rightarrow l^+ \nu$  (11% per lepton)
- $W^+ \rightarrow q qbar$  (33%  $uq$ , 33%  $cq$ )

# SM Top Decays

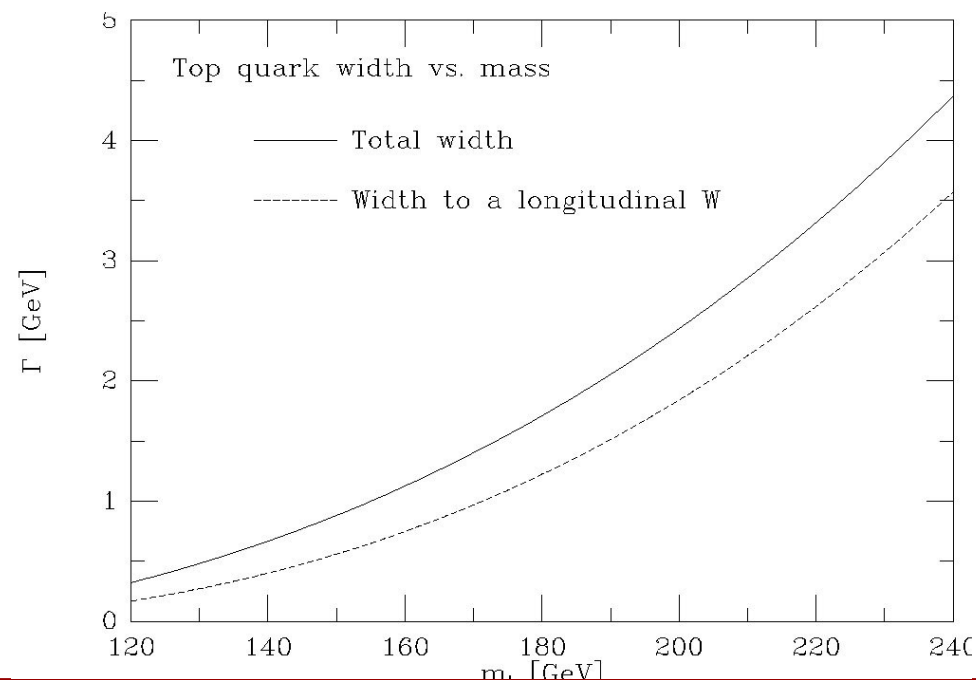
Polarization of W: left-handed (L) or longitudinal (0)

$$\sum |M_L|^2 = \frac{2G_F m_t^4}{\sqrt{2}} |V_{tb}^2| 2x^2(1 - x^2 + y^2)$$

$$\sum |M_0|^2 = \frac{2G_F m_t^4}{\sqrt{2}} |V_{tb}^2| (1 - x^2 - y^2(2 + x^2 - y^2))$$

where  $x=M(W)/m(t)$ ,  $y=m(b)/m(t)$ .

$$\frac{\Gamma(t \rightarrow bW_0)}{\Gamma(t \rightarrow bW)} = \frac{m_t^2}{m_t^2 + 2M_W^2}$$



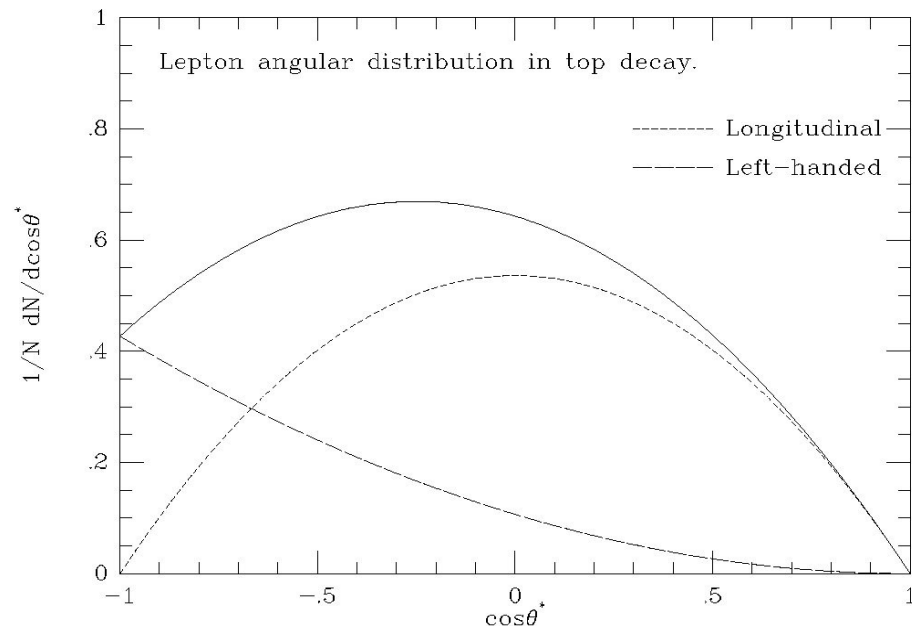
# SM Top Decays

Angular distribution of products:  
depends on W polarization

$$\cos(\theta_e) \approx \frac{4b \cdot e}{m_t^2 - M_W^2}$$

Angle of  $l^+$  in rest frame of W w.r.t. original direction of W

Different distribution: potential  
sign of physics beyond the SM





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# SM Top Decays

Assuming  $m(b), m(W) \rightarrow 0$ , the total width becomes

$$\Gamma_0(t \rightarrow bW) = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}^2| \approx 1.76 \text{ GeV} \left( \frac{m_t}{175 \text{ GeV}} \right)^3$$

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Only one power of the Fermi constant: “Semi-weak” decay rate



# SM Top Decays

Large mass & semi-weak decay rate:

Total decay width **bigger** than typical hadronic scale

Lifetime: order of  $10^{-25}$  sec



# SM Top Decays

Large mass & semi-weak decay rate:

Total decay width **bigger** than typical hadronic scale

Lifetime: order of  $10^{-25}$  sec

=> **top decays before** hadronization

i.e. **no** “toponium” or other hadrons with top-quarks!



# SM Top Decays

**Top quark is as close as we can get to a **free** quark!**



# QCD Corrections

$$\Gamma(t \rightarrow b W) = \Gamma_0 \left[ A^{(0)} + \frac{\alpha_s}{\pi} C_F A^{(1)} + \left( \frac{\alpha_s}{\pi} \right)^2 A^{(2)} + \dots \right]$$



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$A^{(0)}$ : Correction for non-zero  $M_W$

$A^{(1)}$ : Exact result by M. Jezabek & J.H. Kühn in 1988

$A^{(2)}$ : Three different calculations

- A.Czarnecki & K.Melnikov (1998), **assuming**  $M_W \rightarrow 0$
- K.G.Chetyrkin, R.Harlander, T.Seidensticker & M.Steinhauser (1999), Padé approx., **no assumption** for  $M_W$ 
  - Using their notation in the following
- M.Slusarczyk (2004) **analytically**
- Results agree with each other 😊



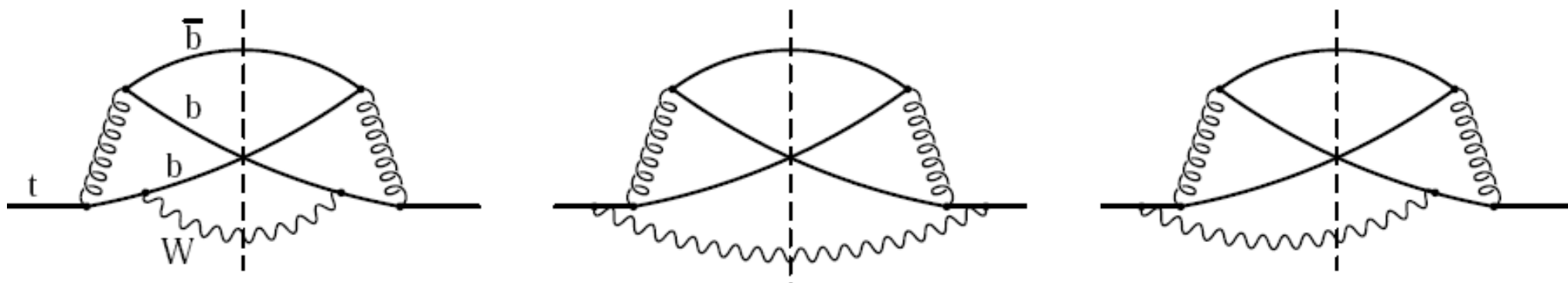
# QCD Corrections - $O(\alpha_s)$

$$A^{(1)} = \frac{5}{4} - \frac{\pi^2}{3} + \frac{3}{2} \frac{M_W^2}{M_t^2} + \frac{M_W^4}{M_t^4} \left( -6 + \pi^2 - \frac{3}{2} \ln \frac{M_t^2}{M_W^2} \right)$$

- M.Jezabek and J.H.Kühn, Nucl Phys **B314**, 1 (1989)
- Calculated “using standard techniques”...
- Exact result

# QCD Corrections - $O(\alpha_s^2)$

- Quite complicated situation
- No exact result available\*
- Calculations involve some non-planar diagrams...



\* : According to a paper from 1999...

# QCD Corrections - $O(\alpha_s^2)$

$$A^{(2)} = C_F^2 A_A^{(2)} + C_A C_F A_{NA}^{(2)} + C_F T n_l A_l^{(2)} + C_F T A_F^{(2)}$$

## Methods used

*Czarnecki & Melnikov:*

- Non-planar diagrams: Expand around  $m_b = m_t/3$  in variable  $\delta = 1 - 3(m_b/m_t)$
- $W$  massless

# QCD Corrections - $O(\alpha_s^2)$

$$A^{(2)} = C_F^2 A_A^{(2)} + C_A C_F A_{NA}^{(2)} + C_F T n_l A_l^{(2)} + C_F T A_F^{(2)}$$

## Methods used

*Chetyrkin et al:*

- Expand each term in  $(M_W/M_t)^2$

$$A_i^{(2)} = A_i^{(2)}|_{M_W=0} + \frac{M_W^2}{M_t^2} A_i^{(2)}|_{M_W^2} + \frac{M_W^4}{M_t^4} A_i^{(2)}|_{M_W^4} + \dots$$

- Calculate the values using Padé approximants

# QCD Corrections - $O(\alpha_s^2)$

$$A^{(2)} = C_F^2 A_A^{(2)} + C_A C_F A_{NA}^{(2)} + C_F T n_l A_l^{(2)} + C_F T A_F^{(2)}$$

## Methods used

*Slusarczyk:*

- Use same expansion
- Calculate coefficients analytically
- Use the optical theorem
  - Need to compute imaginary parts of three-loop self-energy diagrams
  - Various cuts of these diagrams correspond to 2-loop virtual corrections / emission of 1 or 2 real quanta

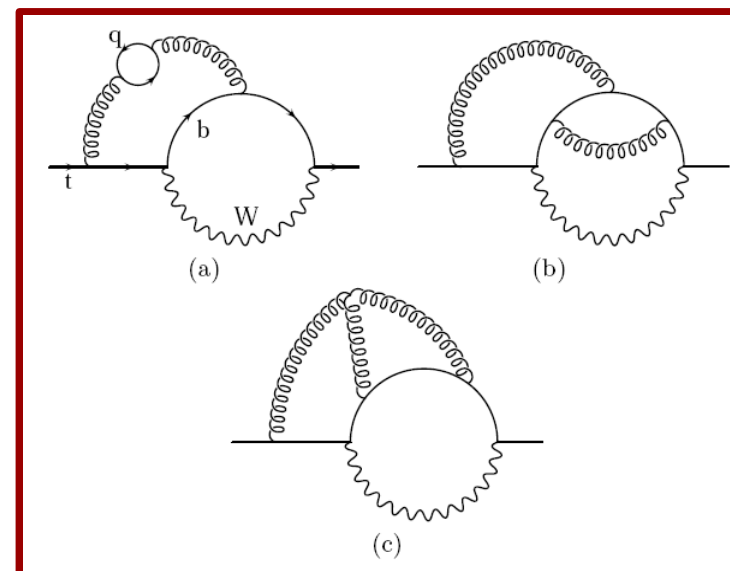
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# QCD Corrections - $O(\alpha_s^2)$

Results of all three methods are quite close

$$A^{(2)} =$$

- -16.7(5)
- -15.6(1.1)
- -15.5(1)

(for  $M_t=175$  GeV,  $M_W=80.4$  GeV,  $\alpha_s(M_t^2)=0.11$ )

# QCD Corrections: Numerically

$$\Gamma(t \rightarrow bW) = \Gamma_0 \left[ A^{(0)} + \frac{\alpha_s}{\pi} C_F A^{(1)} + \left( \frac{\alpha_s}{\pi} \right)^2 A^{(2)} + \dots \right]$$

$A^{(0)}$ : Correction for non-zero  $M_W$       0.8852

$A^{(1)}$ : M. Jezabek & J.H. Kühn      -2.22      ~-10%

$A^{(2)}$ : Three different calculations

– A.Czarnecki & K.Melnikov      -16.7(5)

– K.G.Chetyrkin et al      -15.6(1.1)      ~-2%

– M.Slusarczyk      -15.5(1)

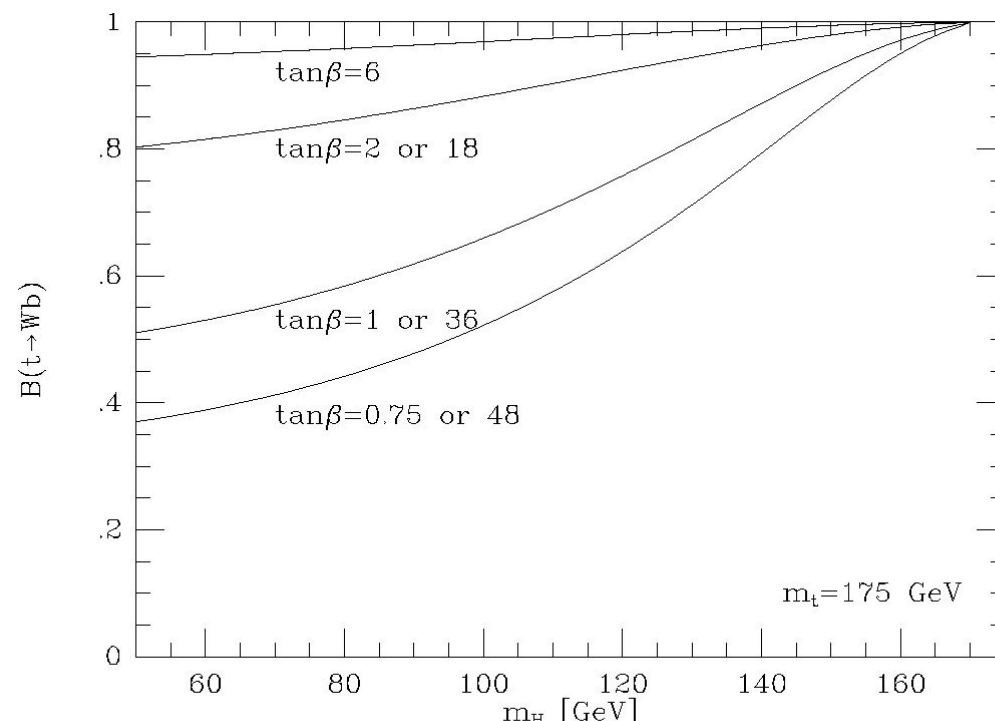
In total:       $\Gamma(t \rightarrow bW) = \sim 0.78 \Gamma_0$

# Bonus: “Unconventional decays”

- MSSM:  $t \rightarrow H^+ b$
- Partial width

$$\Gamma(t \rightarrow Hb) = \frac{G_F m_t^3}{8\Pi\sqrt{2}} |V_{tb}|^2 J\left(\frac{M_H}{m_t}, \frac{m_b}{m_t}, \tan\beta\right)$$

- Has a minimum for  $\tan\beta = \sqrt{m_t/m_b} \sim 6$





# References

- R.K. Ellis, W.J. Stirling, B.R. Webber, *QCD and Collider Physics*, Cambridge University Press 1996
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- A.Czarnecki, K. Melnikov, Nucl Phys **B544**, 520 (1999)
- K.G.Chetyrkin, R.Harlander, T.Seidensticker, M.Steinhauser, Phys Rev D **60**, 114015
- M.Slusarczyk, LLWI 2004 proceedings, 284-288